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Magnetically Self-Bearing Drive System for Ultracentrifugation: Towards 100'000 rpm and $200'000 g$

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Abstract-A novel magnetically double self-bearing drive system for ultracentrifugation in pharmaceutical and chemical industry is presented, built as a prototype and operated. The prototype successfully reached its operating speed of 100'000 rpm. It features an opening stator functionality allowing for fast removal of the ultracentrifuge rotor while enabling future centrifugal acceleration levels of $200'000 g$. The openable stators require a stator encapsulation, provoking unavoidable stator yoke airgaps. Parasitic induced phase-currents were identified, a currentcontrol strategy proposed to diminish them and successfully implemented and operated on the prototype.

Index Terms—self-bearing motor, high-speed, control

I. INTRODUCTION

In biochemistry, fundamental processes are the separation of components of heterogeneous mixtures and purification with ultracentrifuges. Typical ultracentrifuge rotors today are mechanically suspended. Mechanical suspension at high speeds impairs rotor accessibility, requires complicated sealing, causes bearing losses, requires lubrication, generates wear, poses a contamination risk and leads to limited life-time or maintenance needs. These are all limitations of the current technology applied in industry. In [1], the feasibility of a novel system solving these problems was conceptually shown. This article now presents the new fully established working hardware. Furthermore, solutions to problems and challenges related to current control and rotor dynamics which were not yet anticipated in [1] are shown.

In [2], a magnetic levitation based continuous flow ultracentrifuge (UCF) prototype was presented, reaching a centrifugal acceleration C of up to $C = 10^5 g$ and a flow rate of $Q = 0.4$ L/min with a rotational speed higher than 64 krpm. However, the access to the UCF rotor for e.g. cleaning is difficult for the operators due to the self-bearing motor stators fully enclosing it.

This article presents a novel magnetically self-bearing openable drive system, omitting these limitations.

Fig. 1 shows a conceptual illustration thereof. Two selfbearing motors (SBMs), forming a double self-bearing motor, drive the UCF rotor and keep it contact-less in place. As a special feature, a joint between the self-bearing motor modules allows access to the rotor and simple insertion or removal thereof.

Compared to modularly assembled stators during motor manufacturing (e.g. shown in [3]), which are and remain inevitably closed after assembly, the presented self-bearing motors need to be openable by the UCF operators in industry.

This article presents a new hardware UCF drive-system prototype fulfilling the target performance for the industrial UCF application. It furthermore presents solutions to the faced technical challenges in the previously unknown current control strategy of novel openable motors and the rotor-dynamics of the UCF.

In Fig. $2(a)$ the realized prototype of the double self-bearing UCF drive system is shown. It is in the closed ready-to-operate state. Fig. $2(b)$ shows the unique novel ability to unfold in an open state for removal or replacement of the UCF rotor. The indicated system dimensions in Fig. 2 are specified in Tab. I

This article is structured as follows. In Sec. II, two key design challenges are discussed: on the one hand the requirement to have a stator that can be opened and on the other hand the requirement to realize a rotor with bending resonance frequencies above 100 krpm, outside the speed operating range. In Sec. III, the topology choice and its working principle is explained. In Sec. IV, the realized motor prototype is presented, and a fundamental challenge for novel openable

Fig. 1: (a) Conceptual illustration from [1] of the novel magnetically selfbearing drive system for the targeted $100'000$ rpm and $200'000$ g ultracentrifugation. (b) Accessible ultracentrifuge rotor for insertion or removal by operators. (c) Possible placement of the power electronics. (d) Conceptual illustration of the ultracentrifuge rotor.

Fig. 2: (a) Dual self-bearing motor ultracentrifuge drive system prototype in closed ready-to-operate state. (b) Unique novel self-bearing motors being able to unfold into an open state for removal or replacement of the ultracentrifuge rotor by operators. (Indicated system dimensions specified in Tab. I).

motors unveiled. Necessary encapsulation of the stator halves leads to unavoidable yoke air-gaps. The occurrence of resulting flux variations in the stator during operation causes parasitic phase-currents. Therefore, the usage of an L -filter and a voltage feed-forward compensation current control scheme is proposed. In Sec. V, the rotor bending resonances are investigated and sub-critical rotor behavior is shown using 3D-FEM simulation and an acoustic impulse response experiment. The new UCF drive system prototype was successfully commissioned and operated with the proposed current control scheme up to 103'000 rpm. An experimental proof of that high-speed operation and the working of the proposed control method is shown in Sec. VI. The findings are finally summarized in Sec. VII and an outlook is provided.

II. KEY SYSTEM DESIGN CHALLENGES

In $[1]$, the competing system design relationships for a drive system for UCF application were revealed. On the one hand, the rotor diameter of the UCF and its operating speed must be selected such that the desired centrifugal acceleration is achieved $(200'000 g)$ in this article). On the other hand, the drive system must be able to provide necessary mechanical power to the rotor to overcome the air friction losses resulting from the rotor speed, diameter and length. The denser the rotor material, the higher the material load and thus its strength requirement. The rotor material must therefore have a suitable combination of low material density and high strength. In turn, the rotor design must have sufficiently high bending resonance frequencies, determined by the weight and stiffness combination, such that they are not excited during operation.

In the following, the two key challenges are highlighted: on the one hand the requirement for the self-bearing drive system to have an openable stator, on the other hand, the requirement for a sub-critical rotor design for the target speed of 100 krpm.

A. Self-Bearing Motor Opening-Capability

To allow for fast insertion and removal of the UCF rotor by operators in an industrial manufacturing plant, the novel UCF drive-system requires a stator which is split in two separate segments and therefore can be mechanically opened as shown in Fig. $3(a)$.

While open, the stator-halve interface is exposed to the outside production facility/laboratory-environmental conditions. It therefore must be protected from external influences such as contact with process fluids, cleaning agents or unintentional mechanical impacts. Especially the stator core requires special protection. For this purpose, a 0.1 mm thick protective paint coating is envisaged on the exposed stator core separation surfaces. Such a coating behaves magnetically like air, and is therefore referred to in the following as a yoke air-gap in a magnetic sense.

Due to this, in total $0.2 \,\mathrm{mm}$ thin, yoke air-gap in the stator separation plane, harmonic field content is introduced, and the current control needs to be able to counteract resulting induced voltage harmonics.

B. Rotor Dynamics

To avoid excessive vibration or instabilities of the magnetically levitated rotor, the first (and thus lowest) rotor bending resonance frequency, as shown in $Fig. 3(b)$, should be above the maximum frequency of the rotational speed (sub-critical operation). However, higher order synchronous excitations will hit rotor bending modes, and the resulting rotor resonance behavior needs to be stabilized by the control system. The control scheme must furthermore be capable of operating the

novel UCF drive system, despite its new unique properties, including yoke air-gaps in the stator magnetic circuit, which is shown in Sec. IV.

Summarizing, specific design is needed both for the motor as described in Sec. III, Sec. IV, and for the rotor as presented in Sec. V, the validation is in Sec. VI.

III. SLOTLESS DOUBLE SELF-BEARING OPENABLE MOTOR

Self-bearing motors (SBMs) have the unique capability of taking over both the driving and magnetic bearing functionality. This simplifies the drive system architecture, as no separate magnetic bearing units are required.

A wide range of SBM applications and accordingly specialized technical solutions are reported in literature. In [4] an overview over the current state of the technology of SBMs with significant power output is provided. A review of the working principles and topologies of SBMs is published in $[5]$.

The field of SBM research is very active. Current trends include reduction of permanent magnet material usage [6] and new SBM topologies to reduce the complexity of the needed power electronics [7]- [9]. Earlier SBMs had typically separate winding sets for bearing force and torque generation. For combined winding topologies as used in this article, current research addresses their design and operation [10], optimal current utilization strategies [11] and new parallel winding topologies [12]- [15]. General SBM design-guidelines were recently proposed e.g. in [16]- [17], control strategies are active research topics as well, with e.g. [18] investigating the passing through critical speeds and [19] investigating state observers to improve position control. The field of highspeed SBMs experiences continuous attention with recent publications of systems operating at 30 krpm of [20] and [21], at 60 krpm of [22] and at 100 krpm of [23], [24].

However, to the authors' knowledge, the possibility to open high-speed SBM stators in the industrial application by operators, leading to unavoidable yoke air-gaps and their implications on the control, was not reported in literature so far. These aspects are therefore investigated in the following.

A. Topology for UCF Openable SBMs

For the desired UCF application with openable SBM stators, the topology of slotless SBMs with combined windings in toroidal realization as shown in Fig. 4 exhibits several advantages. It allows for the separation of the stator in two separate stator modules.

Furthermore, it leads to a low general and especially rotor loss potential for very high speeds. In contrast, the typical high rotor losses of slotted designs at very high speeds are caused by stronger and higher order harmonics in the airgap field due to the teeth interaction with the field [25], [26]. Slotless SBMs can exhibit an almost sinusoidal air-gap field for the rotor pole-pair number $p = 1$ [27], which is ideal for highspeed applications with low rotor losses. It results in very low harmonic field content, and $p = 1$ leads to the

lowest possible fundamental electrical frequency for a given mechanical speed. This results again in low losses. Therefore for openable UCF SBMs, a slotless stator topology with rotor pole pair number equal to one is very promising.

In SBMs, drive and bearing currents create fields in the same air-gap. The magnetic bearing forces can either be generated with the aid of separate bearing windings, or with the aid of a mathematical superposition of bearing and drive currents in combined windings, as shown in e.g. [28]. The stator field and related forces acting on the rotor are created by superimposed currents of drive and bearing action as shown in e.g. [29]. Combined windings for bearing forces and torque generation lead to equal loading of all coils and full utilization of the winding copper cross-section. If the bearing currents were in separate bearing coil sets, this would not be the case. Therefore, also all semiconductors of the connected power electronics benefit from equal current loading. Additionally, in [27] it was shown that the usage of combined windings lowers the higher order field harmonics induced by the stator drive and bearing currents, additionally reducing the current induced losses at high speeds.

In Fig. $4(d)$ the winding scheme of the resulting UCFs SBM topology is presented. It consists of two three-phase systems a and b with phase connections u_a , v_a , w_a and u_b , v_b , w_b respectively. The two three-phase systems are star-connected at the star-points Y_a and Y_b respectively.

To electrically connect the fixed and the movable stator module in Fig. $4(c)$, a flexible cable connection with a high enough bending radius for repeated bending is envisaged Fig. $4(d)$.

B. Self-Bearing Motor Working Principle

In Fig. 4 the working principle for the UCF SBM torque (a) and force generation (b) is shown. For the torque generation, in Fig. $4(a)$, a stator field with the same pole-pair number as the rotor is generated commonly with all six windings. It is 90° ahead of the rotor's permanent magnet field, which leads to torque generation. To create bearing forces as shown in Fig. 4(b), the same six windings generate a two-pole pair stator field which generates together with the rotor field bearing forces with controllable amplitude and direction.

IV. SELF-BEARING MOTOR WITH YOKE AIR-GAPS

Fig. $5(a)$ shows the stator halves with soft magnetic composite (SMC) cores of one of the realized 100 krpm $48 V_{DC}$ fed 0.25 kW (each) drive system prototypes of the previously discussed topology for UCF application. The main UCF SBM prototype specifications are listed in Tab. II and the used materials of the prototype in Tab. III.

A. Yoke Air-Gap Induced Flux Variation

In Fig. $5(b)$ a 3D-FEM electromagnetic simulation of the magnetic flux density is shown. The yoke air-gaps introduce harmonic field content as shown in Fig. $5(b1)-(b2)$, due to reduced reluctance for $\varphi = \{0, \pi\}$ with no flux crossing over

Fig. 3: Key design challenges: (a) Openable stator for good access to the rotor by UCF operators. Necessary encapsulation of the stator for protection against environmental influences. This leads to unavoidable interruption of the magnetic circuit with corresponding problems for the current-control. (b) Rotor design such that the first bending mode can never be excited by frequencies corresponding to the rotational speed

the yoke air-gaps, and increased reluctance with flux crossing over the yoke air-gap for $\varphi = {\pi/2, -\pi/2}.$

The resulting linked rotor permanent magnet fluxes $\phi_1 - \phi_6$ in the six coils (per turn) are shown in Fig. $5(c)$. The center coils 2 and 5 of the stator limbs experience a slightly higher flux amplitude increased by ϕ_{Δ} . The flux in each phase can therefore be stated as:

$$
\phi_{\mathbf{u}}(t) = (\hat{\phi} + \phi_{\Delta}) \cdot \cos(p\Omega t)
$$

$$
\phi_{\mathbf{v}}(t) = \hat{\phi} \cdot \cos\left(p\Omega t - \frac{2\pi}{3}\right)
$$

$$
\phi_{\mathbf{w}}(t) = \hat{\phi} \cdot \cos\left(p\Omega t + \frac{2\pi}{3}\right), \qquad (1)
$$

where Ω stands for the rotational speed.

Fig. 5(d) shows the dq-transformed fluxes. Both d- and qflux show a second harmonic flux variation with approx. 4% of the d -flux amplitude, as also analytically can be shown by applying the dq -transform on the previous equations, on the one hand for the d-component:

$$
\frac{3}{2}\phi_{d}(t) = \frac{3}{2}\hat{\phi} + \phi_{\Delta} \cdot \cos^{2}(p\Omega t)
$$

=
$$
\frac{3}{2}\hat{\phi} + \frac{1}{2}\phi_{\Delta} + \frac{1}{2}\phi_{\Delta} \cdot \cos(2p\Omega t)
$$
 (2)

and similarly for the q -component:

$$
\frac{3}{2}\phi_{q}(t) = \phi_{\Delta} \cdot \cos(p\Omega t) \cdot \sin(p\Omega t)
$$

$$
= -\frac{1}{2}\phi_{\Delta} \cdot \sin(2p\Omega t)
$$

$$
= \frac{1}{2}\phi_{\Delta} \cdot \cos\left(2p\Omega t - \frac{\pi}{2}\right).
$$

By extracting the time-varying parts of the dq -fluxes $\phi_{d,\Delta}(t)$ and $\phi_{q,\Delta}(t)$:

$$
\phi_{d,\Delta}(t) = \frac{1}{3} \phi_{\Delta} \cdot \cos(2p\Omega t) \n\phi_{q,\Delta}(t) = \frac{1}{3} \phi_{\Delta} \cdot \cos(2p\Omega t - \frac{\pi}{2}),
$$
\n(4)

it is revealed that these time varying components both share the same amplitude, and show the same twofold frequency compared to the rotational speed Ω . They create a 2D-circle in the dq -plane. The same result can be seen by transforming the 3D-FEM obtained fluxes to the dq -plane. It leads to a

rotating flux variation space vector $\underline{\phi}_{\Delta}$, presented in **Fig. 5(e)**.
In the following, the implications, problems and a solution to the discovered flux variation in SBMs with voke air-gaps are unveiled. The presented current-control solution is also realized in the hardware prototype.

B. Parasitic Induced Rotating Current and L-Filter Impact

The separation of the stator into two halves for the stator opening-functionality leads to a parasitic magnetic flux variation. It shows three times the rotational frequency in the stationary frame of reference, and in the moving dq -frame of reference two times the rotational frequency. This leads to an induced voltage space vector \underline{U}_{g} of corresponding frequency. An equivalent circuit of the stator incorporating U_{φ} is shown in Fig. $7(a)$.

The parasitic induced voltage amplitude $|U_{\sigma}|$ can be written as:

$$
|\underline{U}_{g}| = p \cdot N \cdot \phi_{\Delta} \cdot \Omega. \tag{5}
$$

Without further measures, \underline{U}_{g} leads to a parasitic current space
vector \underline{I}_{g} rotating with the same frequency as \underline{U}_{g} as shown (3) in the space vector diagram Fig. 7(b). \underline{U}_{g} therefore leads to

Fig. 4: From [1]: working principle of the slotless SBMs: (a) driving torque, and (b) bearing force generation. (c) Possibility to open the stators of the proposed UCF SBMs enabling access and removal of the UCF rotor. (d) SBM winding scheme of the proposed topology.

parasitic d - and q -currents. In the space vector diagram, the voltage drop across the stator resistance R_S is neglected, as for high efficiency motors $|\underline{U}_{\text{p}}| \gg |\underline{U}_{\text{R}}|$. The parasitic current amplitude $|\underline{I}_{g}|$ calculates as:

$$
|\underline{I}_{g}| = \frac{p \cdot N \cdot \phi_{\Delta} \cdot \Omega}{\sqrt{R_{\rm s}^2 + (3\Omega)^2 \cdot L_{\rm d}^2}}.
$$
 (6)

At low rotational speeds, its amplitude $|\underline{I}_{g}|$ is limited by the stator resistance R_S to an initially linear rise with the speed. At high speeds, it is limited by the motor inductance L_d (rotor is non-salient, i.e. $L_d = L_q$) consisting of an optional phase Lfilter inductance L_f and the stator synchronous inductance L_S , with $L_d = L_F + L_S$.

$$
|I_{\rm g}| = \frac{p \cdot N \cdot \phi_{\Delta} \cdot \Omega}{\sqrt{R_{\rm s}^2 + (3\Omega)^2 \cdot (L_{\rm F} + L_{\rm S})^2}}.\tag{7}
$$

For high speeds, the parasitic current amplitude $|\underline{I}_{g}|$ tends

asymptotically towards the finite value $|\underline{I}_{\varphi,\text{lim}}|$:

$$
|\underline{I}_{g,\text{lim}}| = \frac{p \cdot N \cdot \phi_{\Delta}}{3 \cdot (L_{\text{F}} + L_{\text{S}})}.
$$
 (8)

If the stator winding synchronous inductance L_S is very low, as it is usually the case with high-speed motors, the parasitic current amplitude $|I_{\varphi}|$ can assume large values without additional filter inductances L_F .

Fig. 6 shows for the UCF SBMs of this article the analytically calculated parasitic current amplitude $|I_{\alpha}|$ based on Eq. 7 for a relative flux variation w.r.t. the d-flux amplitude of 4% . It is shown once without L-filter and once with a filter inductance of $L_F = 3 \cdot L_S$. It should be noted, that in the shown 3D-FEM simulation in Fig. $5(b)$, a perfectly manufactured yoke air-gap geometry is assumed. In a real hardware, especially for SMC stator material, the stator edges are imperfect, or need to be rounded slightly, thus increasing the effective yoke air-gap and therefore also the parasitic induced currents I_{σ} , especially for small motors where the non-ideal edge geometry has an even larger influence.

C. Yoke Air-Gap Design

To allow the opening of the SBMs by the industrial operators, the stators are separated in two halves. Necessary protective coating leads unavoidably to at least two yoke airgaps.

From a magnetic standpoint, it seems attractive to place three yoke air-gaps offset by 120° each to achieve a symmetry w.r.t. the three-phase winding. However, the cost and complexity of an industrial system with three stator units, each of them encapsulated and connected by shielded cables leads to the preference of only one separation plane. Additionally, each introduced air-gap increases the unwanted stator reluctance.

The yoke air-gaps should be kept as thin as possible, to reduce adverse effects. However, appropriate durable coating for an industrial environment will always require a certain thickness.

To further reduce the effect of the yoke air-gap between the stator halves, geometric measures are conceivable in the future. On the one hand, the area over which the magnetic flux is transmitted between the stator halves could be increased. On the other hand, if not all coils require the same number of turns for production reasons, the middle coil of each stator segment can have a slightly lower number of turns in order to have the same amount of linked flux in all coils. However, an exact matching, also due to tolerances, seems unlikely. Therefore a current control strategy is always needed.

D. Adverse Effects of Parasitic Induced Rotating Current

The parasitic induced current $I_{\rm g}$ generates a parasitic stator field with pole pair number $p = 1$, rotating with a frequency of $3 \cdot \Omega$ in the stator frame of reference, i.e. an asynchronous field. The self-bearing motors of this article generate bearing fields with number of bearing poles $P_{\text{bng}} = P_{\text{dry}} + 2 = p \cdot 2 +$ $2 = 4$. They are therefore according to [30] of type $P_{\text{bng}} =$ $P_{\text{dry}} \pm 2$. This implies, that the parasitic field cannot generate

Fig. 5: (a) Stator of the realized double self-bearing slotless openable motor prototype. (b) 3D-FEM magnetic field simulation of the prototype motors. Thin yoke air-gaps for durable anti corrosion coating (0.1 mm on each side) in the separation plane. Resulting yoke air-gap introduced harmonic field content (b1)-(b2), due to increased reluctance and flux crossing for $\varphi = 0, \pi$ and reduced reluctance for $\varphi = \pi/2, 3/4\pi$. (c) Resulting linked rotor permanent magnet fluxes $\phi_1 - \phi_6$ in the six coils (per turn). (d) Transformed to dq-quantities, both d- and q-flux show a second harmonic flux variation with approx. 4% of the d-flux amplitude. (e) In dq-plane this leads to a rotating flux variation space vector.

Fig. 6: Analytically calculated parasitic current amplitude $|\underline{I}_{g}|$ for the selfbearing motors of this article given a relative flux variation \tilde{w} .r.t. the d -flux amplitude of 4% once without L-filter and once with a filter inductance of $L_F = 3 \cdot L_S$.

parasitic bearing forces. However parasitic torque ripple \hat{T}_{Ig} of the frequency $2 \cdot \Omega$ is generated. The resulting speed ripple amplitude $\hat{\Omega}_{Ig}$ is kept limited by the rotor inertia J as $\hat{\Omega}$ = T/J and increasingly attenuated with increasing speed Ω by the PI-speed control with low-pass closed-loop characteristic. For the envisaged UCF application, the speed ripple should be as small as possible to avoid re-mixing of the separated media.

 $I_{\rm g}$ furthermore creates additional parasitic motor losses. On the one hand, I_{α} causes additional SMC stator core and rotor losses due to the asynchronous parasitic stator field (with an angular frequency of $3 \cdot \Omega$ in the stator frame of reference). On the other hand, additional copper losses in the windings are generated.

Compared to the rotor PM-field, the parasitic stator field is of less strength, but its threefold frequency of $3 \cdot \Omega$ implies, that SMC core loss terms from the Steinmetz equation [31] with quadratic frequency dependence are multiplied by $9 \cdot \Omega^2$. Furthermore the parasitic field is superimposed to the rotor PM-field and the stator fields for torque and bearing force generation. Thus, the parasitic field is added on top, increasing the loss generation potential, as the core losses are nonlinear with the flux density and sensitive to field bias [32]. A detailed investigation of the complex mechanisms of this loss generation would go beyond the scope of this article and may be subject of future work. This article focuses instead on

the root cause of \underline{I}_{g} and presents and experimentally validates a first approach to mitigate it.

E. Yoke Air-Gap Parasitic Current Feed-forward Control

A voltage compensation by feed-forward control of the parasitic induced voltage U_{σ} on the nominal applied stator voltage $U_{\rm s, nom}$ is proposed as shown in Fig. 7(c) and implemented on the presented prototype. The corresponding control circuit diagram is shown in Fig. 8.

Fig. 7: (a) Stator equivalent electric circuit with parasitic induced voltage $U_{\rm g}$ due to yoke-air-gap flux variation. (b) Space vector diagram with resulting current harmonics $I_{\rm g}$. (c) Proposed feed-forward voltage space vector $U_{\rm ff}$ compensation.

Fig. 8: Proposed current control scheme for the self-bearing openable motors with feed-forward control of the parasitic induced voltage U_{σ} . The power electronics (PE) apply the resulting voltages from drive current-control and position-control to the combined windings of the SBMs.

With the presented voltage feed-forward current control strategy, the adverse effects of the yoke air-gaps can be counteracted.

V. ROTOR DYNAMICS

The conceptual feasibility of an UCF rotor for the drive system presented in this article was shown in [1]. The continuing step from concept to prototype taken in this article requires an investigation of the rotordynamics behavior of the actual design.

Compared to the rotor stiffness, the magnetic bearing stiffness is very small. Therefore, as an analytic approximation for the bending modes, the unsupported free vibration of a Euler-Bernoulli beam can be applied. The first bending mode resonance frequency $f_{res,1}$ of a freely vibrating Euler-Bernoulli beam according to [34] is

$$
f_{\text{res},1} = \frac{22.37}{2\pi L^2} \sqrt{\frac{E \cdot I_x}{\rho \cdot A}},
$$
 (9)

with the length L , Young's modulus E , second moment of area I_x , density ρ and cross-section A. The rotor is modeled as an annulus with inner and outer radii r_1 and r_2 and thickness $t = r_2 - r_1$. With the second moment of area I_x of

$$
I_x - \frac{\pi}{4} (r_2^4 - r_1^4) , \qquad (10)
$$

 $f_{res,1}$ can be rewritten as

$$
f_{\text{res},1} = \frac{22.37}{2\pi L^2} \sqrt{\frac{E}{\rho} \cdot \frac{(r_1 + t)^2 + r_1^2}{4}}.
$$
 (11)

To achieve a high $f_{\text{res},1}$, a short rotor length L, high Young's modulus E, low density ρ , large inner radius r_1 and large thickness t (for given r_1) are beneficial. Therefore these parameters can be varied and the rotor material chosen accordingly, in alignment with the rotor strength and air friction power consumption constraints, as shown in [1]. Knowing these influences, 3D-FEM rotor dynamics simulations can be performed in iterative manner until the desired $f_{res,1}$ is achieved.

Within the scope of this article, two rotor designs were investigated. On the one hand, a test rotor shown in Fig. 9(a1) was designed to verify the drive system with respect to magnetically self-bearing and drive functionality. On the other hand, an UCF rotor prototype displayed in Fig. $9(a2)$ was designed to prove that the rotordynamic UCF requirements can be fulfilled. The rotor dimensions of both rotors are specified in Tab. IV.

For this purpose, 3D-FEM rotor dynamic simulations of the rotor bending modes were performed for both rotors and presented in Campbell diagrams in Fig. $9(b1)$ and Fig. $9(b2)$ respectively. For both rotors, in the speed operating range, the first bending resonance is above the synchronous first-order excitations ($EO = 1$) and thus, it is not excited by them during operation. Excitations of higher order ($EO > 1$) can in principle excite bending modes if they hit their resonance frequency (critical frequencies). However, successful run-up tests in Sec. VI show that these critical frequencies can be passed with the present drive system. Due to the larger rotor diameter compared to the test rotor, the UCF rotor has a higher stiffness and a large reserve with respect to separation of $EO = 1$ and the first bending resonance. To verify the 3D-FEM results, the test rotor was struck with a mechanical impulse while resting stationarily and the acoustic frequency response was measured with a microphone as described in [33] and is shown in Fig. $9(c)$. The provided normalized sound pressure levels p_{rel} are normalized by $p_0 = 20 \mu Pa$ (calibrated at 400 Hz with the sound pressure level measurement unit of [35]). The resonance frequencies determined acoustically and with 3D-FEM agree well, although the deviation is greater for the higher second bending mode. This can be due to the fact, that in the simulation all rotor components were modeled as firmly bonded, resulting in higher stiffness and lower damping than in the prototype with glued components. Thus, the 3D-FEM bending mode simulation was verified acoustically by experiment.

VI. HIGH-SPEED DRIVE-SYSTEM TEST-OPERATION

The UCF drive system presented in this article was successfully commissioned and operated at a speed of 103 krpm with the test-rotor. Thus, the target speed of 100 krpm was reached.

In Sec. IV-B the appearance of parasitic induced currents due to the yoke air-gaps was predicted with the aid of 3D-FEM and analytic derivations. The experimental confirmation thereof is shown in Fig. 10. In Sec. IV-E a feed-forward control compensation was proposed. Its effectiveness is experimentally shown in Fig. 10. Despite the application of L -filters, the parasitic current was found with an amplitude of approx. 0.8 A at 30 krpm to be higher then expected from the FEMsimulations ($> 10\%$ of the nominal UCF operation drive current $\hat{I}_{q,nom}$). This can be a consequence of the non-ideal geometry of the brittle SMC cores with broken off edges, making the interface area between the two core halves smaller then in the ideal simulation geometry. The manufacturing process should therefore be improved in the future. Fig. 10 further shows the achieved reduction of the parasitic currents by approx. 70% with the proposed feed-forward control scheme. Tuning of the feed-forward gain while running levitated at 30 krpm as shown in Fig. 10 allowed to find the optimal gain.

The measured parasitic current amplitude of 0.8 A at 30 krpm leads to a 3D-FEM obtained cross-section average parasitic stator-core flux density component amplitude of 20 mT. The appearance of resulting parasitic current induced loss is proven by measurement. The increase in steady state motor temperature for varying the feed-forward compensation gain from full (100%) to zero and negative compensation (-33%) is shown in Fig. 11 at 20 krpm and 25 krpm. It was found as predicted, that the relative motor temperature, and therefore the parasitic motor loss, grows on the one hand with reduced feed-forward compensation, and on the other hand with increasing speed.

The non-linear growth of the parasitic motor losses with speed as discussed in Sec. IV-D implies the necessity for this 100 krpm self-bearing motors to compensate the parasitic currents.

With the aid of the L -filter and the proposed feed-forward control method it was achieved to reduce the parasitic currents to an acceptable level despite non-ideal stator core manufacturing. With the mechanisms of the generation of the parasitic currents unveiled in this article, the doors are open to test or combine other control methods to counteract the parasitic currents. Work being reported for current regulators in active filters as reviewed in e.g. [36] could be a starting point for future investigations. Figure Fig. 12 shows an oscilloscope screenshot of the rotor angle sensor signal measurement at 103 krpm. Since the rotor has one pole-pair, the fundamental frequency of the shown wave-form of 1.73 kHz corresponds to 103 krpm.

Having reached the target speed, implies that the critical frequencies due to higher order excitations ($EO > 1$) were successfully passed. Fig. 13 shows the radial displacement wave-forms over one rotor-revolution of the rotor angle ϑ . While at 10 krpm some rotor vibrations are visible in Fig. 13(a), at 100 krpm in Fig. 13(b) the rotor is extremely steady. Fig. $13(c)$ shows the average radial displacement radius $|\underline{x}_{r}|_{avg}$ for a speed run-up from 10 – 100 krpm. Crossings of higher order excitations ($\mathbf{EO} > 1$) with the test-rotor bendingmodes resulted in visible displacement radius peaks. However these zones can be successfully passed with the prototype. For unrestricted long-term operation, rotor orbit radii are recommended according to standard [37] to be smaller than $30\% - 40\%$ of the minimal rotor-stator clearance. With 600 μ m clearance, the average rotor orbit radius during speed ramp-up $|\underline{x}_{r}|_{avg}$ in Fig. 13(c) was at resonances always below 12% of the clearance. At the envisaged operating point of 100 krpm it was below 1% . Therefore this recommendation is met and the rotor vibrations sufficiently damped over the whole operating range and the drive system successfully experimentally validated.

VII. CONCLUSION & OUTLOOK

The novel magnetically double self-bearing drive system concept for ultracentrifugation with stator opening functionality for easy rotor removal and performance specifications of $100'000$ rpm and $200'000$ g introduced by the authors in [1] was realized as a prototype. The needed stator encapsulations provoked stator yoke air-gaps were found to lead to parasitic currents if not controlled accordingly. The usage of L -filters and a voltage feed-forward current control-strategy to diminish these effects was proposed and successfully implemented and tested on the newly built prototype. The working of the novel self-bearing openable stator drive-system was experimentally verified. Future work will focus on operating the system in UCF separation-mode, continuing the path towards $200'000 g$.

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Fig. 9: (a1) Test-rotor and (a2) UCF rotor. (b1)-(b2) Corresponding 3D-FEM obtained Campbell diagrams with rotor bending modes. For both rotors, the synchronous excitation with excitation order (EO) equal to one has no intersection with the first bending mode, suggesting the absence of major vibration problems. (c) Acoustic impulse response measurement validation of the test-rotor bending resonance frequencies. The indicated rotor dimensions are specified in Tab. IV.

Fig. 10: Tuning of the parasitic induced voltage feedforward compensation gain conducted at 30 krpm and achieved drastic reduction of the parasitic currents.

Wassmer of Levitronix GmbH for the support during the commissioning process of the novel self-bearing drive-system.

Fig. 11: Relative motor temperature increase for different parasitic induced voltage feed-forward compensation levels (from negative to full compensation) and two speed levels, 20 krpm and 25 krpm.

Fig. 12: Oscilloscope measurement of the rotor angle sensor signal (rotor number of pole pairs equal to one), as a proof of the system operation at 103 krpm.

Fig. 13: Radial rotor displacements of both self-bearing motors over one rig. The natural roof displacements of both sen searing motors over one rotor-revolution (a) at 10 krpm and (b) at 100 krpm. (c) Averaged radial displacement amplitude during a speed ramp-up.

APPENDIX

TABLE I: System Dimensions

Dimension	Variable	Value
System length	$l_{\rm S}$	$340 \,\mathrm{mm}$
System width	$w_{\rm S}$	$260 \,\mathrm{mm}$
System height in closed state	$h_{\rm S,c}$	227 mm
System height in opened state	$h_{\rm S,o}$	$260 \,\mathrm{mm}$

TABLE II: SBM Prototype Specifications

TABLE III: Drive-System Prototype Materials

TABLE IV: Rotor Dimensions

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